

# 0.) Introduction and history

- Description of classical mechanics in terms of Eqs. of motion

→ Newton's second law:

$$\underline{F}(\underline{r}, \underline{p}) = \frac{d}{dt} \underline{p} = \dot{\underline{p}}$$

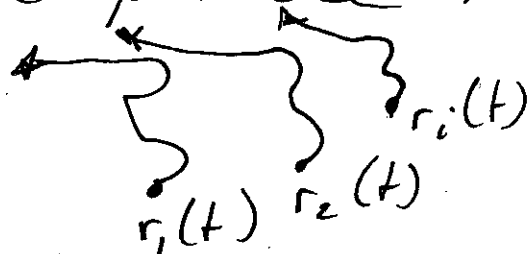
and  $\dot{\underline{r}} = \frac{\underline{p}}{m}$

- Kinematics: description of motion of particles

• Dynamics: origin — " —

- Variables:

- in Euclidean space: 3D  
(one for each particle):



- in configuration space:  $3N - D$

( $N$ : number of particles)

obviously:  $\underline{R}(t) = \{ \underline{r}_1(t), \underline{r}_2(t), \dots \}$

⇒ are equivalent

- phase space :  $2n - 1$

$n$ : number of degrees of freedom,  
 e.g. 3 for mass point in 3D  
 2 for 2 independent points  
 moving on circle  
 etc ...

'2' : for  $r$  and  $p$   
 (coordinate momentum)

→ space of choice for most  
 numerical systems (only first-  
 order derivatives) and  
 quantum mechanics

• History:

Ancient Greeks: very good at math +  
 statics, bad in dynamics/motion

e.g. Pythagoras: concept that complex  
 phenomena can be understood in  
 terms of simpler ones.

- Copernicus/Kepler / Galileo ( $\sim 1500$ )<sup>0-3</sup>

○ Sun at center  
law of planetary motions  
" " gravity

- Descartes ( $\sim 1500$ )  
analytic geometry

- Newton: first full theory of

forces  
behaviors of mass points  
absolute tridimensional space

↑ recognizes that in practice  
both is relative, but absolute  
should exist

(from that later: concept of  
symmetries, relativity)

# I) Survey of the elementary principles

- Review -

## A.) Single particles

velocity:  $\underline{v} = \frac{d}{dt} \underline{r} = \dot{\underline{r}}$

momentum:  $\underline{p} = m \underline{v}$

force:  $\underline{F} = \dot{\underline{p}}$

(for constant force:  $\underline{F} = m \ddot{\underline{r}}$ )

⇒ 1st conservation theorem:  
If no force is applied, the linear momentum  $\underline{p}$  is conserved.

angular momentum:  $\underline{L} = \underline{r} \times \underline{p}$

torque:  $\underline{N} = \underline{r} \times \underline{F}$

or (with

$$\frac{d}{dt} (\underline{r} \times m \underline{v}) = \underline{v} \times m \underline{v} + \underline{r} \times \frac{d}{dt} m \underline{v}$$

$$\underline{N} = \underline{L}$$

$\Rightarrow$  2<sup>nd</sup> conservation theorem:

If total torque is zero, angular momentum is conserved.

(Assume  $m$  constant)

$$\begin{aligned} \text{work: } W_{12} &= \int_1^2 \underline{F} \cdot \underline{ds} \\ &= m \int_{t_1}^{t_2} \frac{d\underline{v}}{dt} \cdot \underline{v} dt = \frac{m}{2} \int_{t_1}^{t_2} \frac{d}{dt} v^2 dt \end{aligned}$$

$$\text{kinetic energy: } = T_2 - T_1 = \frac{m}{2} v_2^2 - \frac{m}{2} v_1^2$$

$\Rightarrow$  if  $W_{12}$  is the same for any path from 1 to 2 (or for

$$\oint \underline{F} \cdot \underline{ds} = 0$$

the force  $\underline{F}$  is conservative.

In this case:  $\int V(r):$

$$\underline{\underline{F}} = - \underline{\nabla} V$$

potential

$$\Rightarrow W_{12} = V_2 - V_1$$

$$\Rightarrow \underline{\underline{T_1 + V_1 = T_2 + V_2}}$$
 energy conservation

$\Rightarrow$  3<sup>rd</sup> conservation theorem: \_\_\_\_\_

If only conservative forces are present, energy  $T + V$  is conserved.

symmetries: Galilean group:

if space is isotropic:  $\underline{r}' = A \underline{r}$  ①

" homogenous:  $\underline{r}' = \underline{r} + \underline{\Delta r}$  ②

if motion is relative:  $\underline{r}' = \underline{r} + \underline{v} \Delta t$  ③

if time is homogenous:  $t' = t + \Delta t$  ④

for all four symmetries: physics does not change between  $\underline{r}', t'$  and  $\underline{r}, t$

$\Rightarrow$  conservation laws