

3.29. Atwood's 3

Define all accelerations positive upwards. By conservation of string, we have $a_2 = -(a_1 + a_3)/2$, because whatever mass disappears above m and $3m$ must appear above $2m$ and be divided evenly between the two segments there. The $F = ma$ equations are

$$T - mg = ma_1, \quad 2T - 2mg = 2ma_2, \quad T - 3mg = 3ma_3. \quad (22)$$

Solving these (the first two quickly give $a_1 = a_2$), along with $a_2 = -(a_1 + a_3)/2$, gives $a_1 = a_2 = g/5$, and $a_3 = -3g/5$.

3.39. Equal distances

We know from Eq. (3.38) that the horizontal distance is $2v_0^2 \sin \theta \cos \theta / g$. The time to the top is $v_0 \sin \theta / g$, so the maximum height is (looking at the ball fall back down to the ground) $gt^2/2 = v_0^2 \sin^2 \theta / 2g$. Equating these results gives $\tan \theta = 4$, so $\theta \approx 76^\circ$.

3.62. Radius of curvature

(a) We have $a = v^2/r \implies r = v^2/a$. At the top, $v = v_0 \cos \theta$ and $a = g$. Therefore, $r = (v_0 \cos \theta)^2 / g$.

(c) The maximum height is the usual $(v_0 \sin \theta)^2 / 2g$. So we want

$$\frac{v_0^2 \cos^2 \theta}{g} = \frac{1}{2} \left(\frac{v_0^2 \sin^2 \theta}{2g} \right) \implies \tan \theta = 2 \implies \theta \approx 63.4^\circ.$$

3.70. Stopping on a cone

The $F = ma$ equation perpendicular to the surface of the cone gives

$$mg \sin \theta - N = (mv^2/R) \cos \theta \implies N = mg \sin \theta - (mv^2/R) \cos \theta. \quad (80)$$

The $F = ma$ equation along the direction of the motion gives $-\mu N = m(dv/dt)$, which yields

$$-\mu dt = \frac{dv}{g \sin \theta - (v^2/R) \cos \theta} \implies -\mu g \sin \theta \int_0^t dt = \int_{v_0}^0 \frac{dv}{1 - \frac{v^2}{gR \tan \theta}}. \quad (81)$$

Letting $u \equiv v/\sqrt{gR \tan \theta}$ gives

$$\begin{aligned} -\mu g \sin \theta \int_0^t dt &= \int_{v_0/\sqrt{gR \tan \theta}}^0 \frac{\sqrt{gR \tan \theta} du}{1 - u^2} \\ \implies t &= -\frac{1}{2\mu} \sqrt{\frac{R}{g \sin \theta \cos \theta}} \ln \left(\frac{1+u}{1-u} \right) \Big|_{v_0/\sqrt{gR \tan \theta}}^0 \\ &= \frac{1}{2\mu} \sqrt{\frac{R}{g \sin \theta \cos \theta}} \ln \left(\frac{\sqrt{gR \tan \theta} + v_0}{\sqrt{gR \tan \theta} - v_0} \right). \end{aligned} \quad (82)$$

Note that if $v_0 = \sqrt{gR \tan \theta}$, then $t = \infty$. This makes sense, because this is the speed for which the string naturally makes an angle of θ with the vertical (as you can show); so the normal force is initially (and hence always) equal to zero. Also, if $\theta \rightarrow \pi/2$ (more precisely, if $v_0/\sqrt{gR \tan \theta} \ll 1$), then to lowest order the argument of the log is $1 + 2v_0/\sqrt{gR \tan \theta}$. So the log is essentially equal to $2v_0/\sqrt{gR \tan \theta}$. We then obtain $t \approx v_0/(\mu g \sin \theta) \approx v_0/(\mu g)$, which makes sense because the acceleration on flat ground is simply $a = -\mu g$. (Or more generally for $\theta \neq \pi/2$, if v_0 is very small, the normal force is essentially equal to $mg \sin \theta$, so the acceleration is $a = -\mu g \sin \theta$.)