

Lecture 03

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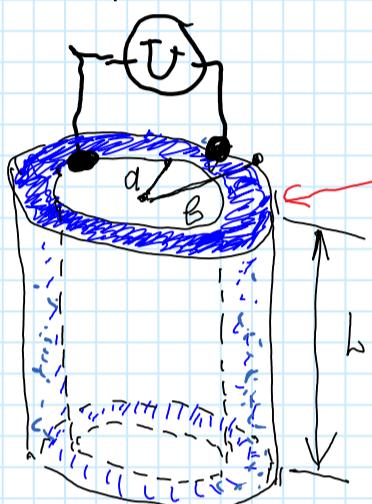
Examples of EM Problems and their Solutions

Ohm's Law: $\boxed{U=RI \text{ or } \vec{J} = \sigma \cdot \vec{E}}$ (σ - conductivity)

EM Induction Law: $\boxed{\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \text{ or } \vec{E} = -\nabla \phi - \frac{\partial \vec{A}}{\partial t}}$

Problem 1: Ohm's law and resistance

Example: Calculate the resistance R of the segment of two coaxial cylinders. The voltage U is applied between A and B surfaces.



Ohm's Law: $\boxed{R = U/I}$

I -current; U - voltage

We can calculate the electric current I and difference of potentials $\phi(r_a) - \phi(r_b) = U$, if the electric charge is distributed on

the surface of the inner cylinder of the radius "a". We can assign the linear charge density λ for the inner cylinder charge distribution.

Electric field for the region $a < r < b$:

$$E = \frac{\lambda}{2\pi\epsilon_0 r}$$

and the density of the electric current J :

$$J = \sigma E = \frac{\sigma \lambda}{2\pi\epsilon_0 r}$$

The difference of potentials U :

$$U = \int_a^b E dr = \frac{\lambda}{2\pi\epsilon_0} \int_a^b \frac{dr}{r} = \frac{\lambda}{2\pi\epsilon_0} \ln b/a;$$

The electric current I :

$$I = J \cdot 2\pi r L = \frac{\lambda \cdot b}{2\pi\epsilon_0} \cdot 2\pi L = \frac{\lambda b L}{2\pi\epsilon_0}$$

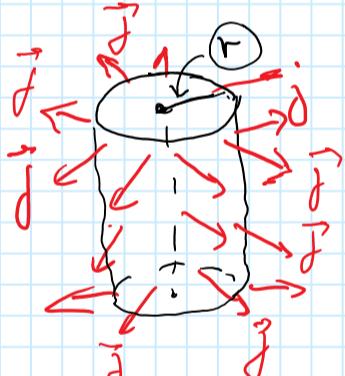
The resistance R :

$$R = \frac{U}{I} = \frac{\lambda b \ln b/a \cdot \epsilon_0}{\lambda \cdot b L} \Rightarrow R = \frac{\ln b/a}{2\pi\epsilon_0 L}$$

recap:

Gauss' Theorem

$$\oint \vec{E} d\vec{s} = E \cdot 2\pi r L = \frac{q}{\epsilon_0} = \frac{\lambda L}{\epsilon_0}$$



Problem 2:

Find the electric field E induced by the cylindrically symmetric magnetic field:

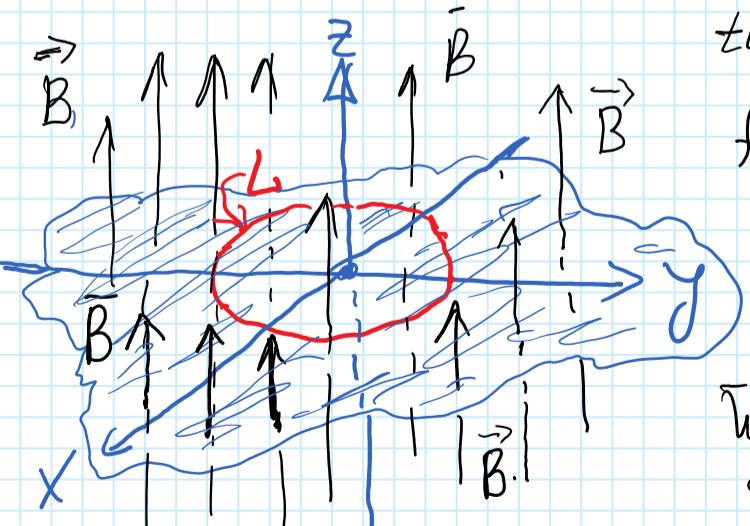
$$\vec{B}(r, t) = \hat{e}_z \cdot B_0 \cdot \cos \omega t$$

the unit vector

of the z -axis

$B_0, \omega \equiv \text{constants}$

We will discuss two methods of solution.

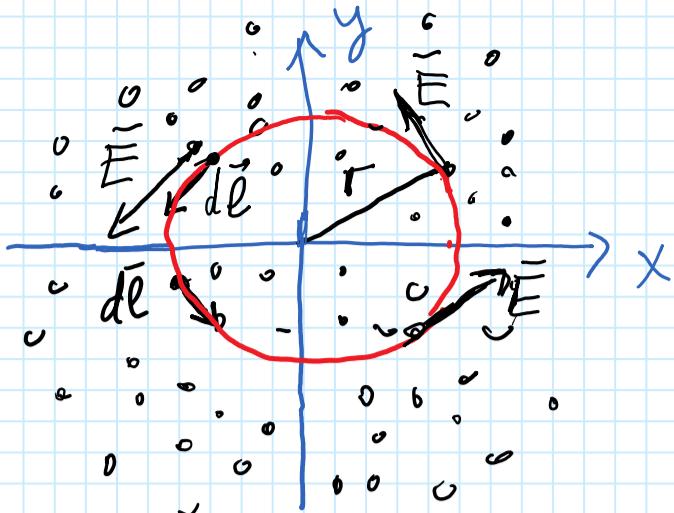


Method 1

Calculation of emf: $\oint \vec{E} d\vec{l} = -\frac{\partial}{\partial t} \int_L \vec{B} d\vec{s}$

The magnetic field B is uniform and we can use a circular loop "L" to calculate emf (the red circle in the figure)

- 2 -



$$\oint \vec{E} \cdot d\vec{l} = E(r) \cdot 2\pi r = - \frac{d}{dt} \int_S \vec{B} \cdot d\vec{s} = -\pi r^2 \frac{dB}{dt}$$

due to cylindrical symmetry

$\vec{B} \parallel \hat{e}_z$
and $d\vec{s} = ds \cdot \hat{e}_z$

$$E(r) \cdot 2\pi r = -\pi r^2 \frac{dB}{dt} \Rightarrow \boxed{\vec{E}(r) = -\frac{r}{2} \frac{dB}{dt}}$$

$$\frac{dB}{dt} = \frac{d}{dt} (B_0 \cos \omega t) = -\omega B_0 \sin \omega t$$

$$\boxed{\vec{E}(r) = \frac{\omega r B_0}{2} \sin \omega t}$$

Method 2: Direct calculation of the vector \vec{E} from the equation describing the Faraday's Law of EM induction:

$$\vec{E} = -\vec{\nabla} \phi - \frac{\partial \vec{A}}{\partial t}$$

$$\phi = 0 \quad \vec{E} = -\frac{\partial \vec{A}}{\partial t}$$

$$\vec{A} = \frac{1}{2} (\vec{B} \times \vec{r}) = \frac{1}{2} B r \cdot \hat{e}_\phi$$

$$\vec{E} = +\frac{1}{2} r \frac{\partial B}{\partial t} \hat{e}_\phi \Rightarrow$$

$$\boxed{\vec{E} = +\frac{\omega r B_0}{2} \sin \omega t \cdot \hat{e}_\phi}$$

In our Problem $\phi = 0$. The vector potential of the cylindrically symmetric uniform magnetic field can be written as

$$\boxed{\vec{A} = \frac{1}{2} (\vec{B} \times \vec{r})}. \text{ Let us check it:}$$

$$\vec{B} = \vec{\nabla} \times \vec{A} = \frac{1}{2} \vec{\nabla} \times (\vec{B} \times \vec{r}) = \frac{1}{2} [\vec{B}(\vec{\nabla} \cdot \vec{r}) - (\vec{B} \cdot \vec{\nabla})\vec{r}]$$

common vector algebra

$$\vec{B}(\vec{a} \vec{c}) - \vec{c}(\vec{a} \vec{b}) = \vec{B}(\vec{a} \vec{c}) - (\vec{b} \vec{a})\vec{c}$$

$$\vec{B} \vec{a} \vec{b} \vec{c} = \vec{B} \frac{d}{dz} \vec{r} = \vec{B}$$

$$\vec{\nabla} \cdot \vec{r} = \frac{\partial x}{\partial x} + \frac{\partial y}{\partial y} + \frac{\partial z}{\partial z} = 3;$$

$$\vec{B} = \vec{\nabla} \times \vec{A} = \frac{1}{2} (3\vec{B} - \vec{B}) = \vec{B} \quad \text{All is correct!}$$

For the cylindrical coordinates (r, ϕ, z) :

$$\vec{A} = \frac{1}{2} (\vec{B} \times \vec{r}) = \frac{1}{2} B r \cdot \hat{e}_\phi$$