

Preliminary Exam: Electrodynamics, Thursday January 14, 2016. 9:00-12:00

Answer a total of any **THREE** out of the four questions.

For your answers you can use either the blue books or individual sheets of paper.

If you use the blue books, put the solution to each problem in a separate book.

If you use the sheets of paper, use different sets of sheets for each problem and sequentially number each page of each set.

Be sure to put your name on each book and on each sheet of paper that you submit.

If you submit solutions to more than three problems, only the first three problems as listed on the exam will be graded.

Some Possibly Useful Information

Inside covers of Jackson Classical Electrodynamics enclosed.

Problem 1

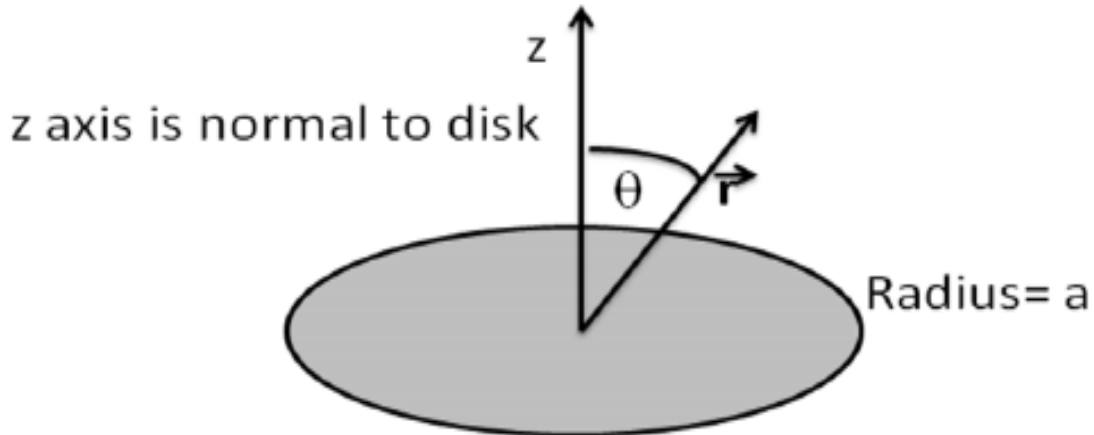
An amount of charge q is uniformly spread on the surface of a disk of radius a .

- (a) Calculate the potential as a function of z on the axis of symmetry indicated in the figure.
- (b) Calculate the potential at any point \mathbf{r} with $r > a$. Express the result in angular harmonics ($P_n(\cos \theta)$), where θ is the angle between the z -axis and \mathbf{r} . Hint: Solve Laplace's equation in spherical coordinates, not cylindrical coordinates. The solution is a series of the form $\sum A_n f_n(r) P_n(\cos \theta)$ summed over $n = 0, 1, 2, 3, \dots$, where $f_n(r)$ is a function of r and the A_n are constants. You may find the following expansion useful.

$$\frac{1}{|\mathbf{r} - \mathbf{r}'|} = \frac{1}{r} \sum_{\ell=0}^{\infty} \left(\frac{r'}{r}\right)^{\ell} P_{\ell}(\cos \gamma)$$

where γ is the angle between \mathbf{r} and \mathbf{r}' .

- (c) Use the answer of part (a) to determine the constants in part (b).



Problem 2

Consider a plane polarized electromagnetic wave propagating in a medium of dielectric constant ϵ , permeability μ , and conductivity σ . The electric field is given by

$$\mathbf{E} = \mathbf{E}_0 e^{i\omega t - i\omega n \mathbf{k} \cdot \mathbf{r}/c}$$

where \mathbf{k} is the wavevector, \mathbf{r} is the position vector, c is the speed of light in vacuum, ω is the frequency, t is the time, and n is the complex refractive index. Using Maxwell's equations (with $\mathbf{D} = \epsilon \mathbf{E}$, $\mathbf{B} = \mu \mathbf{H}$, and $\mathbf{J} = \sigma \mathbf{E}$), obtain an expression for the complex index of refraction n in terms of ω , ϵ , μ , and σ .

Problem 3

A point (perfect) electric dipole \mathbf{p} is located near a grounded conducting sphere of radius R . The radius-vector connecting the center of the sphere and the dipole is \mathbf{r}_0 where $r_0 > R$.

- (a) Construct the electrostatic image of a dipole using the expression for the Green's function in the presence of a conducting sphere

$$G(\mathbf{r}, \mathbf{r}_0) = \frac{1}{|\mathbf{r} - \mathbf{r}_0|} - \frac{Rr_0}{|r_0^2\mathbf{r} - R^2\mathbf{r}_0|}$$

where \mathbf{r}_0 is the location of a unit charge source.

- (b) Determine the electrostatic potential $\phi(\mathbf{r})$ outside the sphere, if the dipole vector \mathbf{p} is perpendicular to the vector \mathbf{r}_0 .
- (c) Determine the potential $\phi(\mathbf{r})$, if the dipole vector \mathbf{p} is oriented along the vector \mathbf{r}_0

Problem 4

The density $\mathbf{J}(\mathbf{r})$ of a **stationary** electric current propagating along the z -direction is given by the vector

$$\mathbf{J}(\mathbf{r}) = \mathbf{e}_z J_0 \cos(\mathbf{k} \cdot \mathbf{r})$$

where \mathbf{r} is the radius-vector, J_0 is a positive constant, and the vector $\mathbf{k} = k_x \mathbf{e}_x + k_y \mathbf{e}_y$ is orthogonal to the unit vector \mathbf{e}_z of the z -axis.

- (a) What are the Poisson equations for the three projections of the vector potential $\mathbf{A}(\mathbf{r})$ of the **induced magnetostatic** field?
- (b) Determine the vector-potential $\mathbf{A}(\mathbf{r})$ induced by the current \mathbf{J} , by solving the Poisson equation for the vector potential. The boundary conditions for the vector potential and magnetic field are $\mathbf{A}(\mathbf{r} = 0) = 0$ and $\mathbf{B}(\mathbf{r} = 0) = 0$.
- (c) Find the locations where the density of magnetic energy $w_B(\mathbf{r})$ is minimal.

Vector Formulas

$$\begin{aligned}
 \mathbf{a} \cdot (\mathbf{b} \times \mathbf{c}) &= \mathbf{b} \cdot (\mathbf{c} \times \mathbf{a}) = \mathbf{c} \cdot (\mathbf{a} \times \mathbf{b}) \\
 \mathbf{a} \times (\mathbf{b} \times \mathbf{c}) &= (\mathbf{a} \cdot \mathbf{c})\mathbf{b} - (\mathbf{a} \cdot \mathbf{b})\mathbf{c} \\
 (\mathbf{a} \times \mathbf{b}) \cdot (\mathbf{c} \times \mathbf{d}) &= (\mathbf{a} \cdot \mathbf{c})(\mathbf{b} \cdot \mathbf{d}) - (\mathbf{a} \cdot \mathbf{d})(\mathbf{b} \cdot \mathbf{c}) \\
 \nabla \times \nabla \psi &= 0 \\
 \nabla \cdot (\nabla \times \mathbf{a}) &= 0 \\
 \nabla \times (\nabla \times \mathbf{a}) &= \nabla(\nabla \cdot \mathbf{a}) - \nabla^2 \mathbf{a} \\
 \nabla \cdot (\psi \mathbf{a}) &= \mathbf{a} \cdot \nabla \psi + \psi \nabla \cdot \mathbf{a} \\
 \nabla \times (\psi \mathbf{a}) &= \nabla \psi \times \mathbf{a} + \psi \nabla \times \mathbf{a} \\
 \nabla(\mathbf{a} \cdot \mathbf{b}) &= (\mathbf{a} \cdot \nabla)\mathbf{b} + (\mathbf{b} \cdot \nabla)\mathbf{a} + \mathbf{a} \times (\nabla \times \mathbf{b}) + \mathbf{b} \times (\nabla \times \mathbf{a}) \\
 \nabla \cdot (\mathbf{a} \times \mathbf{b}) &= \mathbf{b} \cdot (\nabla \times \mathbf{a}) - \mathbf{a} \cdot (\nabla \times \mathbf{b}) \\
 \nabla \times (\mathbf{a} \times \mathbf{b}) &= \mathbf{a}(\nabla \cdot \mathbf{b}) - \mathbf{b}(\nabla \cdot \mathbf{a}) + (\mathbf{b} \cdot \nabla)\mathbf{a} - (\mathbf{a} \cdot \nabla)\mathbf{b}
 \end{aligned}$$

If \mathbf{x} is the coordinate of a point with respect to some origin, with magnitude $r = |\mathbf{x}|$, $\mathbf{n} = \mathbf{x}/r$ is a unit radial vector, and $f(r)$ is a well-behaved function of r , then

$$\begin{aligned}
 \nabla \cdot \mathbf{x} &= 3 & \nabla \times \mathbf{x} &= 0 \\
 \nabla \cdot [\mathbf{n}f(r)] &= \frac{2}{r}f + \frac{\partial f}{\partial r} & \nabla \times [\mathbf{n}f(r)] &= 0 \\
 (\mathbf{a} \cdot \nabla)\mathbf{n}f(r) &= \frac{f(r)}{r} [\mathbf{a} - \mathbf{n}(\mathbf{a} \cdot \mathbf{n})] + \mathbf{n}(\mathbf{a} \cdot \mathbf{n}) \frac{\partial f}{\partial r} \\
 \nabla(\mathbf{x} \cdot \mathbf{a}) &= \mathbf{a} + \mathbf{x}(\nabla \cdot \mathbf{a}) + i(\mathbf{L} \times \mathbf{a})
 \end{aligned}$$

where $\mathbf{L} = \frac{1}{i}(\mathbf{x} \times \nabla)$ is the angular-momentum operator.

Theorems from Vector Calculus

In the following ϕ , ψ , and \mathbf{A} are well-behaved scalar or vector functions, V is a three-dimensional volume with volume element d^3x , S is a closed two-dimensional surface bounding V , with area element da and unit outward normal \mathbf{n} at da .

$$\int_V \nabla \cdot \mathbf{A} d^3x = \int_S \mathbf{A} \cdot \mathbf{n} da \quad (\text{Divergence theorem})$$

$$\int_V \nabla \psi d^3x = \int_S \psi \mathbf{n} da$$

$$\int_V \nabla \times \mathbf{A} d^3x = \int_S \mathbf{n} \times \mathbf{A} da$$

$$\int_V (\phi \nabla^2 \psi + \nabla \phi \cdot \nabla \psi) d^3x = \int_S \phi \mathbf{n} \cdot \nabla \psi da \quad (\text{Green's first identity})$$

$$\int_V (\phi \nabla^2 \psi - \psi \nabla^2 \phi) d^3x = \int_S (\phi \nabla \psi - \psi \nabla \phi) \cdot \mathbf{n} da \quad (\text{Green's theorem})$$

In the following S is an open surface and C is the contour bounding it, with line element $d\mathbf{l}$. The normal \mathbf{n} to S is defined by the right-hand-screw rule in relation to the sense of the line integral around C .

$$\int_S (\nabla \times \mathbf{A}) \cdot \mathbf{n} da = \oint_C \mathbf{A} \cdot d\mathbf{l} \quad (\text{Stokes's theorem})$$

$$\int_S \mathbf{n} \times \nabla \psi da = \oint_C \psi d\mathbf{l}$$

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Explicit Forms of Vector Operations

Let $\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3$ be orthogonal unit vectors associated with the coordinate directions specified in the headings on the left, and A_1, A_2, A_3 be the corresponding components of \mathbf{A} . Then

Cartesian
($x_1, x_2, x_3 = x, y, z$)

$$\begin{aligned}\nabla\psi &= \mathbf{e}_1 \frac{\partial\psi}{\partial x_1} + \mathbf{e}_2 \frac{\partial\psi}{\partial x_2} + \mathbf{e}_3 \frac{\partial\psi}{\partial x_3} \\ \nabla \cdot \mathbf{A} &= \frac{\partial A_1}{\partial x_1} + \frac{\partial A_2}{\partial x_2} + \frac{\partial A_3}{\partial x_3} \\ \nabla \times \mathbf{A} &= \mathbf{e}_1 \left(\frac{\partial A_3}{\partial x_2} - \frac{\partial A_2}{\partial x_3} \right) + \mathbf{e}_2 \left(\frac{\partial A_1}{\partial x_3} - \frac{\partial A_3}{\partial x_1} \right) + \mathbf{e}_3 \left(\frac{\partial A_2}{\partial x_1} - \frac{\partial A_1}{\partial x_2} \right) \\ \nabla^2\psi &= \frac{\partial^2\psi}{\partial x_1^2} + \frac{\partial^2\psi}{\partial x_2^2} + \frac{\partial^2\psi}{\partial x_3^2}\end{aligned}$$

Cylindrical
(ρ, ϕ, z)

$$\begin{aligned}\nabla\psi &= \mathbf{e}_1 \frac{\partial\psi}{\partial\rho} + \mathbf{e}_2 \frac{1}{\rho} \frac{\partial\psi}{\partial\phi} + \mathbf{e}_3 \frac{\partial\psi}{\partial z} \\ \nabla \cdot \mathbf{A} &= \frac{1}{\rho} \frac{\partial}{\partial\rho} (\rho A_1) + \frac{1}{\rho} \frac{\partial A_2}{\partial\phi} + \frac{\partial A_3}{\partial z} \\ \nabla \times \mathbf{A} &= \mathbf{e}_1 \left(\frac{1}{\rho} \frac{\partial A_3}{\partial\phi} - \frac{\partial A_2}{\partial z} \right) + \mathbf{e}_2 \left(\frac{\partial A_1}{\partial z} - \frac{\partial A_3}{\partial\rho} \right) + \mathbf{e}_3 \frac{1}{\rho} \left(\frac{\partial}{\partial\rho} (\rho A_2) - \frac{\partial A_1}{\partial\phi} \right) \\ \nabla^2\psi &= \frac{1}{\rho} \frac{\partial}{\partial\rho} \left(\rho \frac{\partial\psi}{\partial\rho} \right) + \frac{1}{\rho^2} \frac{\partial^2\psi}{\partial\phi^2} + \frac{\partial^2\psi}{\partial z^2}\end{aligned}$$

Spherical
(r, θ, ϕ)

$$\begin{aligned}\nabla\psi &= \mathbf{e}_1 \frac{\partial\psi}{\partial r} + \mathbf{e}_2 \frac{1}{r} \frac{\partial\psi}{\partial\theta} + \mathbf{e}_3 \frac{1}{r \sin\theta} \frac{\partial\psi}{\partial\phi} \\ \nabla \cdot \mathbf{A} &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_1) + \frac{1}{r \sin\theta} \frac{\partial}{\partial\theta} (\sin\theta A_2) + \frac{1}{r \sin\theta} \frac{\partial A_3}{\partial\phi} \\ \nabla \times \mathbf{A} &= \mathbf{e}_1 \frac{1}{r \sin\theta} \left[\frac{\partial}{\partial\theta} (\sin\theta A_3) - \frac{\partial A_2}{\partial\phi} \right] \\ &\quad + \mathbf{e}_2 \left[\frac{1}{r \sin\theta} \frac{\partial A_1}{\partial\phi} - \frac{1}{r} \frac{\partial}{\partial r} (r A_3) \right] + \mathbf{e}_3 \frac{1}{r} \left[\frac{\partial}{\partial r} (r A_2) - \frac{\partial A_1}{\partial\theta} \right] \\ \nabla^2\psi &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial\psi}{\partial r} \right) + \frac{1}{r^2 \sin\theta} \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\partial\psi}{\partial\theta} \right) + \frac{1}{r^2 \sin^2\theta} \frac{\partial^2\psi}{\partial\phi^2} \\ &\quad \left[\text{Note that } \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial\psi}{\partial r} \right) \equiv \frac{1}{r} \frac{\partial^2}{\partial r^2} (r\psi). \right]\end{aligned}$$