

Oral Examination



# Observation of CP Violation in Kaon Decays

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# Outline

1. Theory
  - i. Review of discrete symmetries
  - ii. Violation of symmetries and the CPT Theorem
  - iii. The Kaon System
    - a. Mixing and eigenstates of CP
    - b. Decay modes and regeneration
2. Experiment: Cronin, Fitch et. al.
  - i. Setup
  - ii. Analysis
  - iii. Results



# Discrete Symmetries

The discrete transformations discussed here have eigenvalues

$$\hat{O} |\psi\rangle = \lambda |\psi\rangle \quad \lambda = +1, -1$$

where the operator stands in for Charge Conjugation (C), Parity (P), and Time Reversal (T)

What are their eigenstates?

If  $[\hat{O}, \hat{H}] = 0$ , that is if  $\hat{O}$  has the same eigenstates as the Hamiltonian, then these energy eigenstates are said to have definite states of symmetry.



# CPT Theorem

*A local, Lorentz invariant quantum field theory with a Hermitian Hamiltonian must respect CPT symmetry.*

- first appeared in the work of Julian Schwinger, then proven more explicitly by Lüders, Pauli and Bell.
- stands on solid ground theoretically and experimentally

*Implications:* individual violations of permutations of C, P and T must cancel. Thus, violation of CP would require violation of T, which would mean that

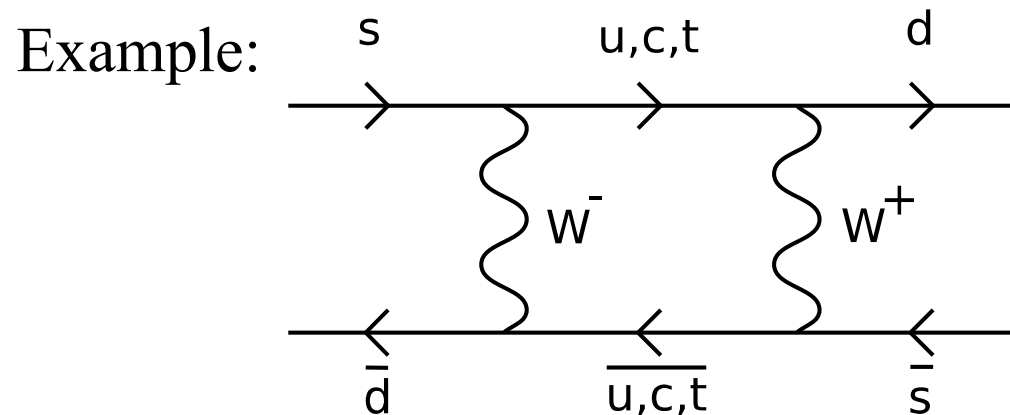
- time has a preferred direction on the fundamental scale.
- there is a clue to the matter-antimatter imbalance (the two are otherwise CP-symmetric)

# The Kaon System

Neutral Kaon Particles:  $K^0 = d\bar{s}$ ;  $\bar{K}^0 = \bar{d}s$

- Neutral particle with a distinct (opposite strangeness) antiparticle
- Common decay products (e.g.  $2\pi$ )

*Consequence:* A neutral Kaon can oscillate into its antiparticle!



These must not be eigenstates of the full Hamiltonian!

# The Kaon System

## Mixing Formalism:

Evidently, the strong interaction Hamiltonian\*:

$$H_{strong} = \begin{pmatrix} M & 0 \\ 0 & M \end{pmatrix} \quad \text{eigenstates:} \quad K^0 = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \bar{K}^0 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

acquires off-diagonal “mixing terms” due to the weak interaction:

$$H = \begin{pmatrix} M & V \\ V & M \end{pmatrix} \quad \begin{array}{l} \text{eigenstates:} \quad K_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} \quad K_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} \\ \text{eigenvalues:} \quad E_1 = M + V \quad E_2 = M - V \end{array}$$

Time evolution

introduces oscillation:

$$K_1 e^{-\frac{i}{\hbar}(M+V)t} + K_2 e^{-\frac{i}{\hbar}(M-V)t} = \sqrt{2} e^{-\frac{i}{\hbar}Mt} \begin{pmatrix} \cos \frac{V}{\hbar} t \\ i \sin \frac{V}{\hbar} t \end{pmatrix}$$

$K_1$  and  $K_2$  (imaginary) decay rates are added on the diagonal

\* Rest frame assumed to avoid extra contributions to the energy.

# The Kaon System

## Neutral Kaons as states of CP Transformation

*Problem:* Kaons are not good states of CP:  $CP(K^0) = -\bar{K}^0$

...but the eigenstates of the new Hamiltonian are:

$$K_1 = \frac{1}{\sqrt{2}}(K^0 - \bar{K}^0) \Rightarrow CP(K_1) = \frac{1}{\sqrt{2}}(-\bar{K}^0 + K^0) = K_1 \quad \langle CP \rangle = +1$$
$$K_2 = \frac{1}{\sqrt{2}}(K^0 + \bar{K}^0) \Rightarrow CP(K_2) = \frac{-1}{\sqrt{2}}(\bar{K}^0 + K^0) = -K_2 \quad \langle CP \rangle = -1$$

Success? CP and the Hamiltonian have simultaneous eigenstates – CP must be conserved, i.e. symmetry states maintained:

$$K_1 \rightarrow 2\pi \quad \langle CP \rangle_{2\pi} = +1$$
$$K_2 \rightarrow 3\pi \quad \langle CP \rangle_{3\pi} = -1$$

Is this true or can we find:  $K_2 \rightarrow 2\pi \quad \langle CP \rangle: -1 \rightarrow +1$

# The Kaon System

## Experimental Perspective

	$\tau$ (s)	Main decay modes	$\Gamma_i / \Gamma$	Experimental use
$K_1$	$\sim 10^{-10}$	$\pi^+ \pi^-$ $\pi^0 \pi^0$	69.2% 30.7%	← useful for calibration, conveniently short lifetime
$K_2$	$\sim 10^{-8}$	$\pi^+ l^- \nu_l$ or conj. ( $K_{l3}$ ) $3\pi^0$ $\pi^+ \pi^- \pi^0$	67.6% 19.6% 12.6%	interesting potential source of CP violation; can <i>regenerate</i> $K_1$

## Regeneration

$$\dot{K}_2 = \frac{1}{\sqrt{2}} \left( K^0 + \bar{K}^0 \right) \left. \begin{array}{l} \begin{array}{l} \xrightarrow{\quad} \bar{K}^0 + p \rightarrow \Lambda^0 + \pi^+ \\ \xrightarrow{\quad} K^0 + p \rightarrow \Lambda^0 + \underbrace{K^0 + K^+} \end{array} \\ \text{strong interactions:} \\ \text{must conserve strangeness} \end{array} \right\}$$

leave little free energy – unlikely!

$K^0$  remains, so  $K_1$  is back! (in superposition with  $K_2$ )



# Experimental Setup

BROOKHAVEN  
NATIONAL LABORATORY

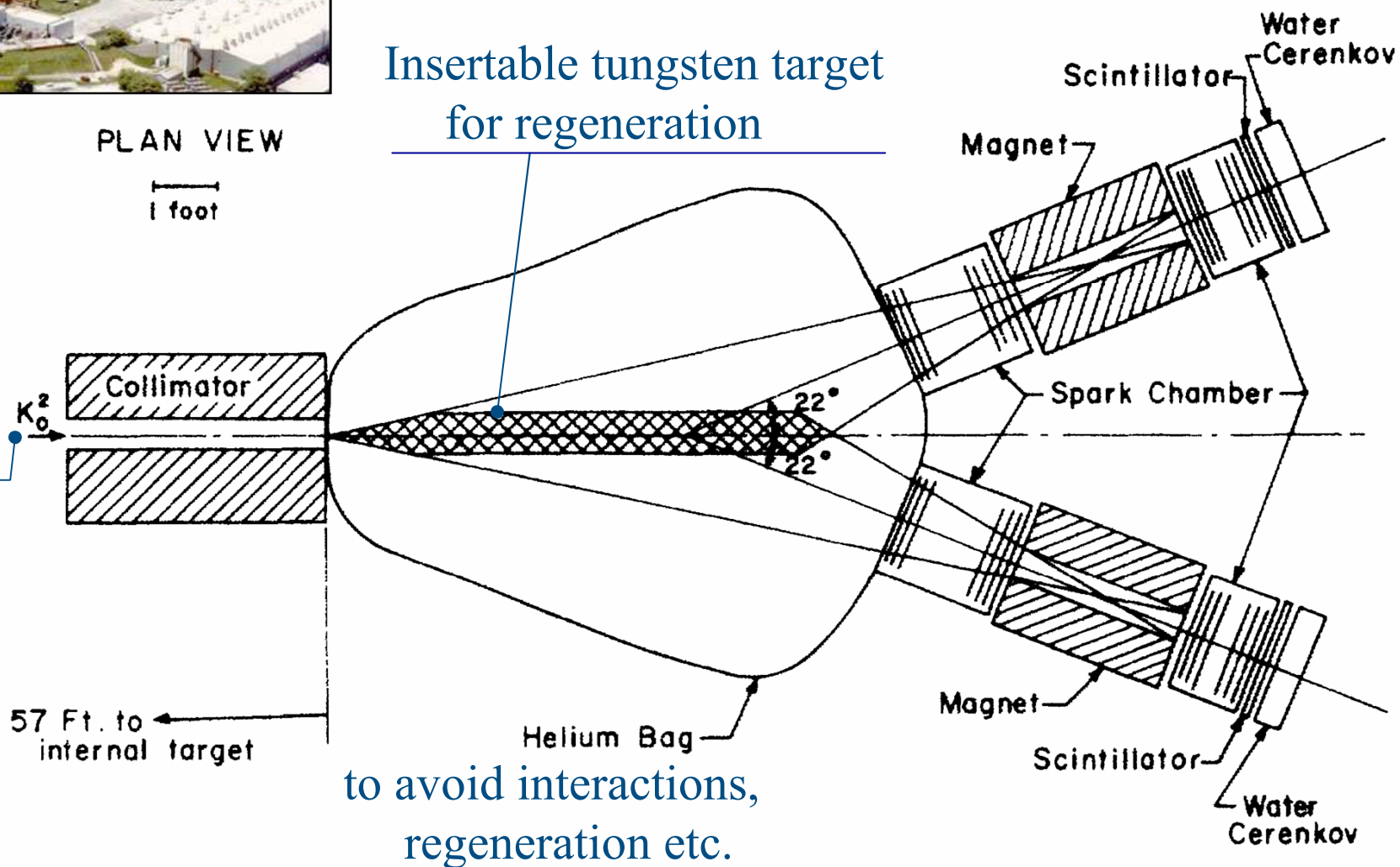


AGS

PLAN VIEW

1 foot

$K_1$  decayed  
away by this  
point





# Data Analysis

$2\pi$  decay filtering method:

- both particles are captured: invariant mass of  $K^0$  expected
- forward direction ( $\theta = 0$ ) for the vector sum of the two momenta

Not so for other possible (3-body) decays –  $K_{e3}$ ,  $K_{\mu3}$ ,  $K_{\pi3}$ : decay products are lost. Result:

- invariant mass is undercounted
- $\theta \neq 0$

Approach to calibration and measurement

Regenerate  $K_1$  and measure  $\theta$  and  $m$  distributions of  $2\pi$  decay and compare with those of  $K_2$  if such decays are found.

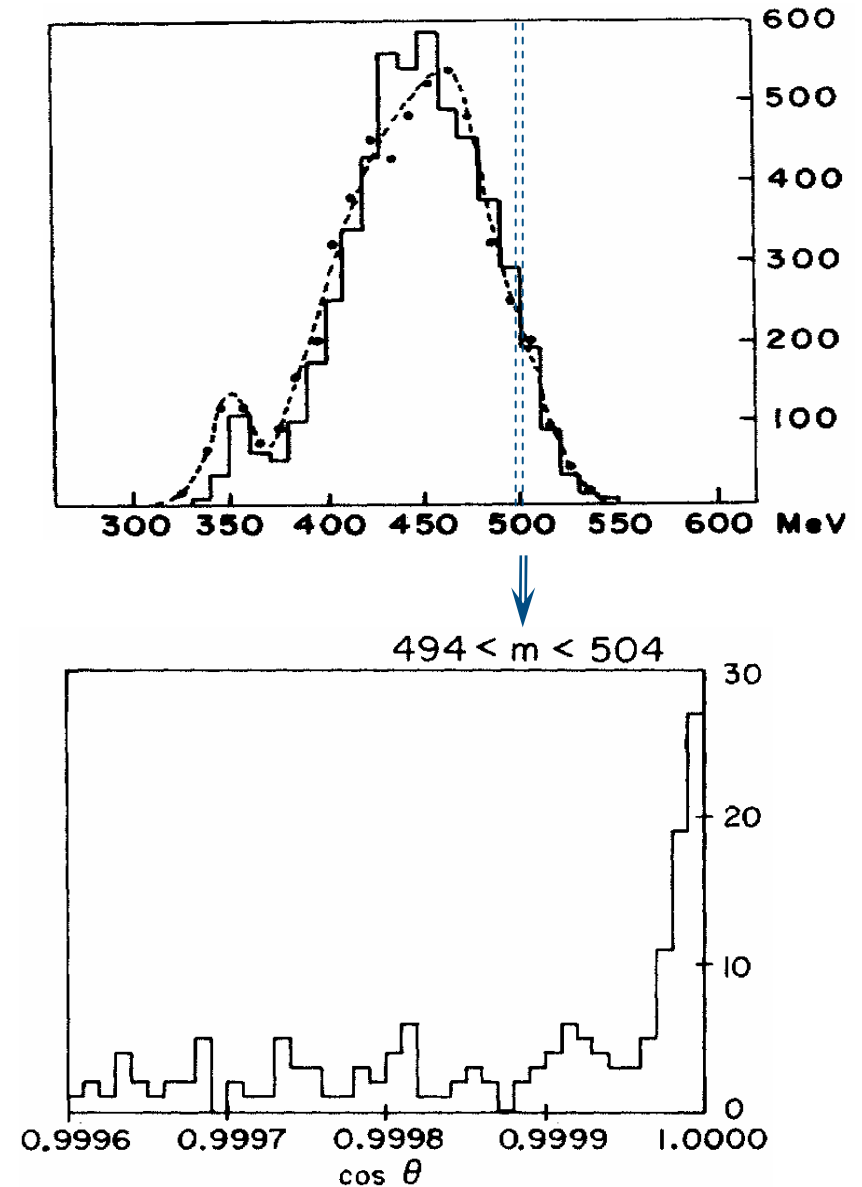
# Data Analysis

The result of “mass undercounting”:  
mass spectrum spreads and shifts  
below the  $K^0$  mass.

Cutting on  $K^0$  mass and looking for a  
forward peak in the  $\cos \theta$  distribution  
(sign of 2-body decay)...

**$2\pi$  decay invariant mass and angle  
distributions are the same as those  
from regenerated  $K_1$**

	inv. mass (MeV)	peak angle (mrad)
$K_1$	$498.1 \pm 0.4$	$3.4 \pm 0.3$
$K_2$	$499.1 \pm 0.8$	$4.0 \pm 0.7$



# Results

So, having subtracted the background as shown and taken into account relative detection efficiencies, there were found  $45 \pm 9$  CP-violating  $\pi^+\pi^-$  decays out of a total of 22700 events. This corresponds to a branching ratio of  $0.20 \pm 0.04 \%$ .

Reported:

VOLUME 13, NUMBER 4

PHYSICAL REVIEW LETTERS

27 JULY 1964

## EVIDENCE FOR THE $2\pi$ DECAY OF THE $K_2^0$ MESON\*†

J. H. Christenson, J. W. Cronin,† V. L. Fitch,‡ and R. Turlay§

Princeton University, Princeton, New Jersey

(Received 10 July 1964)

This Letter reports the results of experimental studies designed to search for the  $2\pi$  decay of the  $K_2^0$  meson. Several previous experiments have

The analysis program computed the vector momentum of each charged particle observed in the decay and the invariant mass,  $m^*$ , assuming



# Results

Evidently, the short and long-lived particles (i.e. energy eigenstates having distinct decay rates) previously thought to be eigenstates of CP are in fact:

$$K_S^0 \approx K_1^0 + \varepsilon K_2^0$$

$$K_L^0 \approx K_2^0 + \varepsilon K_1^0$$

where  $K_1$  and  $K_2$  are the pure eigenstates of CP and  $\varepsilon$  is the degree of violation. Calculated in the analysis of the original experiment:

$$|\varepsilon| = 2.3 \times 10^{-3}$$



# Summary

The presented results lead to the following conclusions:

- the Weak interaction slightly violates CP symmetry
- by the CPT theorem, it violates T symmetry as well – a preferred direction on the elementary particle scale!
- a small (and not yet satisfactory) degree of CP violation has been verified in the theory of matter-antimatter imbalance.