

Holographic Pomeron and Schwinger Mechanism

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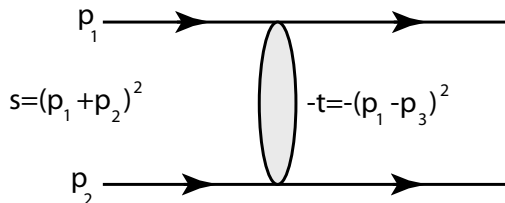
April 30, 2012

GB, D. Kharzeev, H.U. Yee & I. Zahed
arXiv:1202.0831, Phys.Rev. D**85** (2012) 105005

Outline

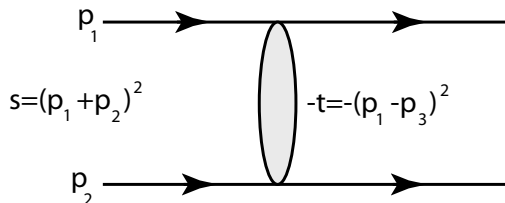
- ▶ Motivation and basics
- ▶ Calculation of the scattering amplitude
- ▶ Schwinger string production
- ▶ Conclusions

Motivation



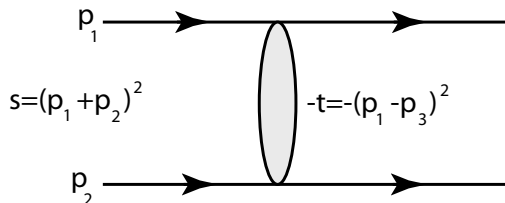
- Understand near forward, high energy ($s \gg -t$) hadronic scattering amplitudes

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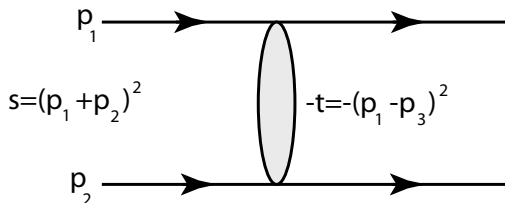
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- ▶ Small momentum transfer $\rightarrow$ non-perturbative
- ▶ Large impact parameter $\rightarrow$ confinement
- ▶ Total cross section: $\sigma_{tot}(s) = s^{-1} \text{Im} [\mathcal{T}(s, t = 0)]$

Regge theory

Analyticity and crossing symmetry of the S matrix

$$\mathcal{T}(s, t) \approx \beta(t) \frac{s^{\alpha(t)} \pm (-s)^{\alpha(t)}}{\sin(\pi\alpha(t))}$$

Singularities of $\mathcal{T}$ in the complex angular momentum plane



Exchange of families of states in t-channel (“Reggeons”)



Regge trajectory: $\alpha(t)$

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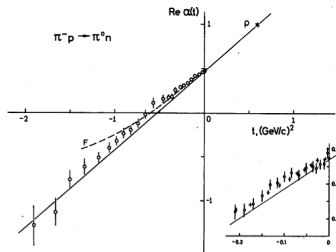
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Exchange of families of states in t-channel (“Reggeons”)



Regge trajectory: $\alpha(t) = \alpha(0) + \alpha' t$



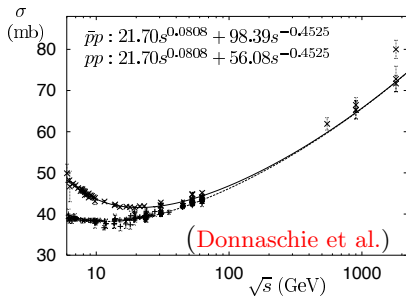
(Apel et al. 79)

Pomeron

$$\sigma_{tot} \propto s^{\alpha(0)-1}$$

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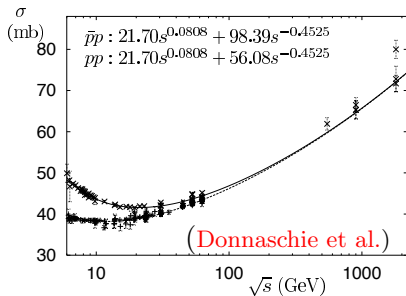
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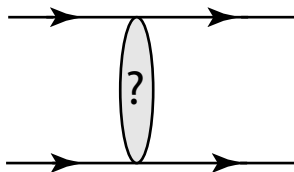
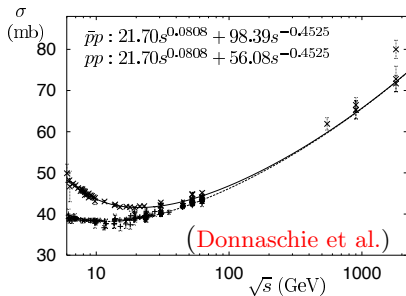
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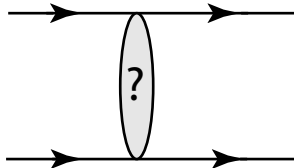
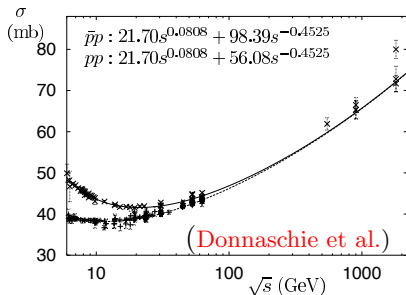


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Pomeron: Regge trajectory with $\alpha_{\mathbb{P}}(0) \approx 1$

- ▶ (**Donnachie, Landshoff**), fit to experiment, $\alpha(0) \approx 1.08$
- ▶ (**Low-Nussinov**), two gluon exchange, $\alpha(0) = 1$
- ▶ (**BFKL**), gluon ladder, (cylinder topology at large N_c), $\alpha(0) \approx 1.3$
- ▶ ...



Idea

Calculate dipole-dipole amplitude through gauge/string duality

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Dipole $\sim$ Wilson loop

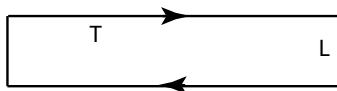
Wilson lines

$$\mathbf{W}(\mathcal{C}) = \frac{1}{N_c} \text{Tr} \left[\mathbf{P}_c \exp \left(ig \int_{\mathcal{C}} A_{\mu}(x) dx^{\mu} \right) \right]$$

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Static Wilson loop:



$\langle \mathbf{W} \rangle \approx e^{-V(L)T}$ $V(L)$: potential between 2 static charges

- ▶ Abelian gauge field: $V(L) \propto 1/L \Rightarrow$ “Perimeter law”
- ▶ Confining potential: $V(L) = \sigma_T L \Rightarrow$ “Area law”

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- ▶ Recoil of scatterers is negligible $\rightarrow$ follow fixed trajectories
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gluon $\sim$ color dipole at large N_c

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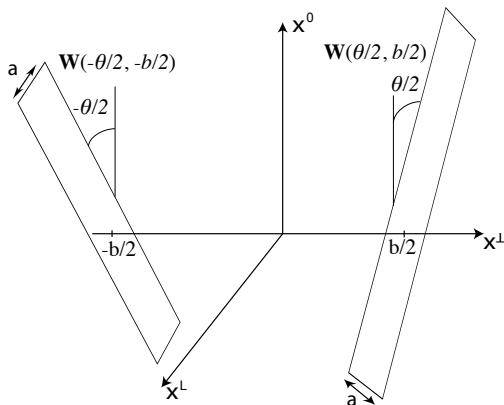
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high energy hadron $\sim$ gluon cloud $\sim$ bunch of dipoles
gluon $\sim$ color dipole at large N_c

Dipole-dipole kinematics

$$D_1(p_1) + D_2(p_2) \rightarrow D_1(k_1) + D_2(k_2)$$



$$p_1/m = (\cos(\theta/2), -\sin(\theta/2), 0^\perp)$$

$$p_2/m = (\cos(\theta/2), \sin(\theta/2), 0^\perp)$$

$$q = (0, 0, q^\perp)$$

$$b = (0, 0, b^\perp)$$

$$\theta \rightarrow -i\chi$$

$$\chi \approx \log(s/2m^2) \gg 1$$

(Meggiolaro, Giordano,
Janik, Peschanski,
Nowak, Shuryak, Zahed...)

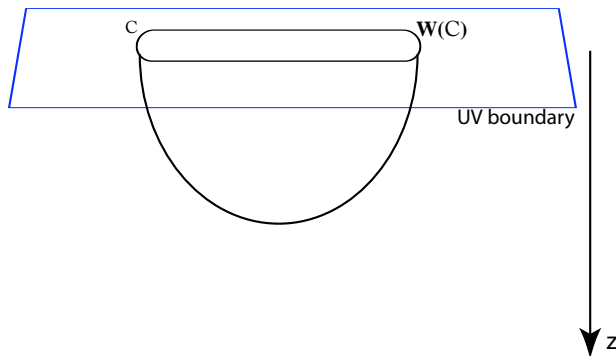
(Nachtmann, '91)

$$\frac{1}{-2is} \mathcal{T}(\theta, q) \approx \int d^2b e^{iq^\perp \cdot b} \langle (\mathbf{W}(\theta/2, b/2) \mathbf{W}(-\theta/2, -b/2) - \mathbf{1}) \rangle$$

Wilson lines in holography

$$\langle \mathbf{W}\mathbf{W} \rangle = Z_{string}[\partial\mathcal{B} = \mathcal{C}] \quad (\text{Maldacena, '98})$$

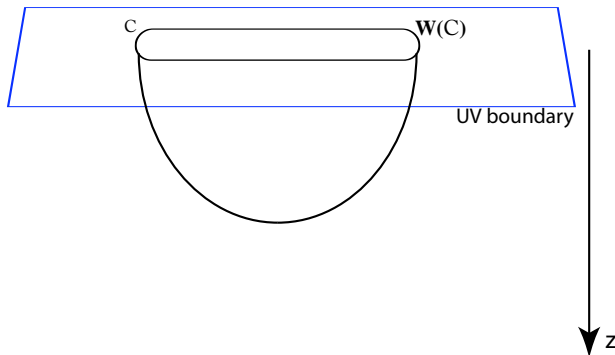
Deconfined Geometry



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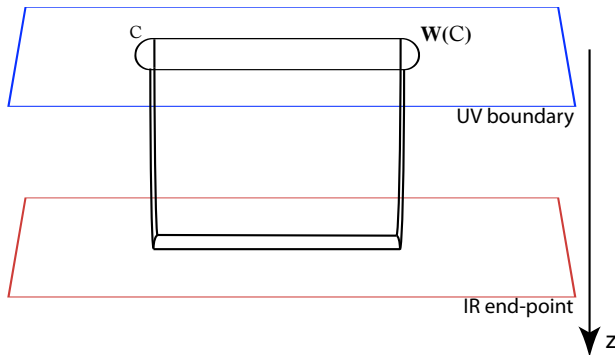
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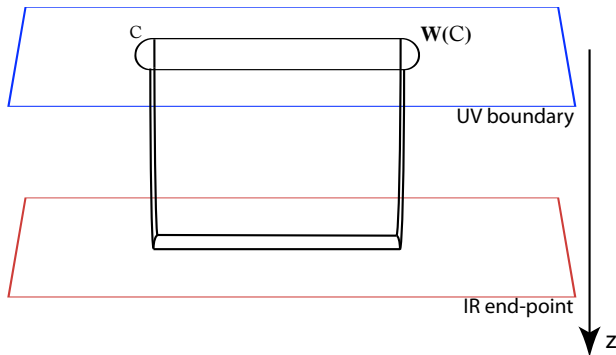
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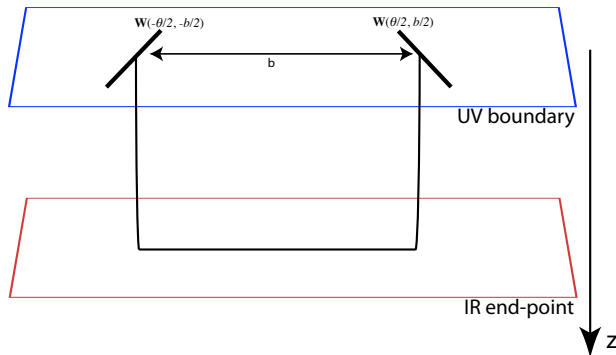
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Confining Geometry (Witten, '98)



String partition function

1-loop (cylinder) partition function:

$$\langle \mathbf{W}\mathbf{W} \rangle = g_s^2 \int_0^\infty \frac{dT}{2T} \int_{\mathbf{T}} d[x] e^{-\mathbf{S}[x] + \text{ghosts}} = g_s^2 \int_0^\infty \frac{dT}{2T} \mathbf{K}(T)$$

$$\mathbf{S} = \frac{\sigma_T}{2} \int_0^T d\tau \int_0^1 d\sigma (\dot{x}^\mu \dot{x}_\mu + x'^\mu x'_\mu)$$

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Boundary conditions:

$$\text{time :} \quad x^\mu(T, \sigma) = x^\mu(0, \sigma)$$

$$\begin{aligned} \text{space :} \quad & \cos(\theta/2) x^1(\tau, 0) + \sin(\theta/2) x^0(\tau, 0) = 0 \\ & \cos(\theta/2) x^1(\tau, 1) - \sin(\theta/2) x^0(\tau, 1) = 0 \end{aligned}$$

String partition function

$$\mathbf{K}(T) = \frac{\pi \sigma_T a^2}{\sinh(\theta T/2)} e^{-\sigma_T b^2 T/2} \eta^{-D_\perp}(iT/2) \prod_{n=1}^{\infty} \prod_{s=\pm} \frac{\sinh(\pi n T/2)}{\sinh(\pi(n + s\theta/\pi)T/2)}$$

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$\frac{1}{\sinh(\theta T/2)} \prod_{n=1}^{\infty} \prod_{s=\pm} \frac{\sinh(\pi n T/2)}{\sinh(\pi(n + s\theta/\pi)T/2)}$: longitudinal modes

Back to Minkowski space: $\theta \rightarrow -i\chi$

$$\langle \mathbf{W}\mathbf{W} \rangle = \frac{ia^2 g_s^2}{4\alpha'} \int_0^\infty \frac{dT}{T} \frac{1}{\sin(\chi T/2)} e^{-\sigma_T b^2 T/2} \eta^{-D_\perp}(iT/2) \Pi(..)$$

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Multi-wrapping: $k \equiv \text{“N-ality”}$, $k > 1$: N_c suppressed

Leading pole:

$$\langle \mathbf{W}\mathbf{W} \rangle = -a^2 \pi g_s^2 (4\pi\alpha')^{D_\perp-1} \left(\frac{1}{2\pi\alpha'\chi} \right)^{D_\perp/2} e^{-b^2/2\alpha'\chi + D_\perp\chi/12}$$

Diffusion term

$$\langle \mathbf{W}\mathbf{W} \rangle = -a^2 \pi g_s^2 (4\pi\alpha')^{D_\perp-1} \left(\frac{1}{2\pi\alpha' \chi} \right)^{D_\perp/2} e^{-b^2/2\alpha' \chi + D_\perp \chi/12}$$

Diffusion in transverse space (Gribov '75)

$$\mathbf{K}(\chi, b) = \left(\frac{1}{2\pi\alpha' \chi} \right)^{D_\perp/2} e^{-b^2/2\alpha' \chi}$$

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Rapidity $\chi \leftrightarrow$ “diffusion time” $\alpha' \leftrightarrow$ diffusion constant

Lüscher term and the intercept

$$\langle \mathbf{W}\mathbf{W} \rangle = -a^2 \pi g_s^2 (4\pi\alpha')^{D_\perp-1} \left(\frac{1}{2\pi\alpha' \chi} \right)^{D_\perp/2} e^{-b^2/2\alpha' \chi + D_\perp \chi/12}$$

Lüscher term (Casimir energy):

$$\frac{D_\perp \chi}{12} = \frac{\pi D_\perp}{6} \frac{b}{T} = b \cdot V(T) \quad , \quad T = \frac{2\pi b}{\chi}$$

$$e^{\frac{D_\perp \chi}{12}} = s^{D_\perp/12} \text{ (intercept)}$$

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Fermions?

$$V(T)_F \sim D_\perp \sqrt{\frac{M_{KK}}{T}} e^{-2M_{KK}T} \quad (\text{Sonnenschein et al., '00})$$

Scattering amplitude and cross section

$$\begin{aligned}\mathcal{T}(s,t) &= -2is \int d^2b e^{iq_\perp b} \langle (\mathbf{W}\mathbf{W} - 1) \rangle \\ &= ia^2 \pi^2 g_s^2 \left(\frac{\pi}{\ln s} \right)^{D_\perp/2-1} s^{\alpha(t)}\end{aligned}$$

Regge trajectory: $\alpha(t) = 1 + \frac{D_\perp}{12} + \frac{\alpha'}{2}t$

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Total cross section: $\sigma_{tot} = a^2 \pi^2 g_s^2 \left(\frac{\pi}{\ln s} \right)^{D_\perp/2-1} s^{D_\perp/12}$

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Total cross section: $\sigma_{tot} = a^2 \pi^2 g_s^2 \left(\frac{\pi}{\ln s} \right)^{D_\perp/2-1} s^{D_\perp/12}$

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Scattering amplitude and cross section

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- Sum over non-interacting Pomeron exchanges:

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- Saturates the Froissart bound

Schwinger mechanism and Regge trajectory

$$\text{scattering amplitude} \sim \langle \mathbf{W} \mathbf{W} \rangle \sim e^{-\frac{b^2}{2\alpha'\chi}} = e^{-\frac{\pi m_s^2}{\sigma_T \chi}}$$

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Intermezzo: Schwinger pair production

Worldline formalism (**Dunne,Schubert**):

$$\Gamma[A] = \int_0^\infty \frac{dT}{T} e^{-m^2 T} \int_{PBC} d[x] \exp \left[- \int_0^T d\tau \left(\frac{\dot{x}^2}{4} + i A \cdot \dot{x} \right) \right]$$

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$$\Gamma_s = g_s^2 \int_0^\infty \frac{dT}{2T} \int_{PBC} d[x] \exp \left[- \int_0^T d\tau \int_0^1 d\sigma (\dot{x}^2 + x'^2) \right]$$

Schwinger string production

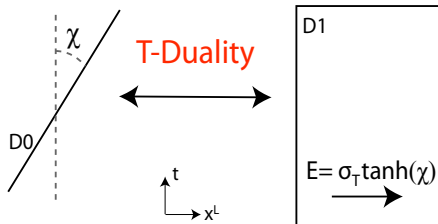
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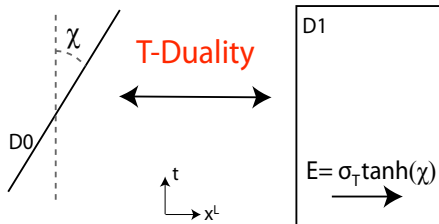
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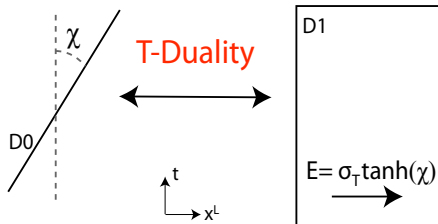


$$\text{twisted b.c.} \Leftrightarrow \left. \frac{E}{2} \int d\tau (\tilde{x}^1 \partial_\tau x^0 - x^0 \partial_\tau \tilde{x}^1) \right|_{\sigma=0,1}$$

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“Worldsheet instantons” (**Schubert**):

$$x^0 = R(\sigma) \cos(2\pi k u) \ , \ \tilde{x}^1 = R(\sigma) \sin(2\pi k u) \ , \ x^T = b(\sigma - 1/2)$$

$$R(\sigma) = \frac{b}{\chi} \cosh(\chi(\sigma - 1/2)) = \frac{m}{\sigma_T \chi} \cosh(\chi(\sigma - 1/2))$$

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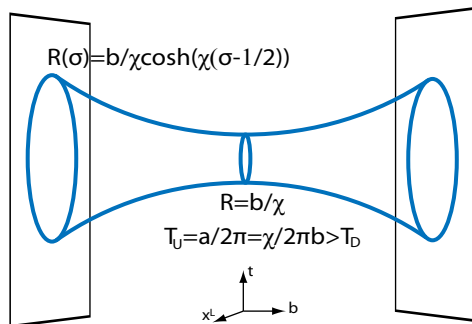
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Schwinger mechanism and Regge trajectories

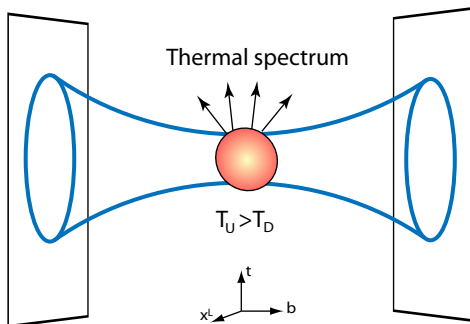
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Electric field $\leftrightarrow$ Acceleration $\leftrightarrow$ Unruh effect



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Conclusions:

- ▶ nonperturbative approach to soft pomeron (slope and intercept)
- ▶ inelasticity & Regge behavior $\Leftrightarrow$ string creation à la Schwinger
- ▶ fireball in the center of the collision with thermal radiation
- ▶ relative rapidity $\chi \Leftrightarrow$ electric field $E_L = \sigma_T \tanh \chi$